Structural Stability for a Class of Nonlinear Wave Equations

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ABSTRACT

In this paper we discuss the structural stability of an initial value problem defined for the equation $u_t - u_{txx} + \alpha u u_x = \beta u_x u_{xx} + u u_{xxx}$

where α , β are constants, $x \in \mathbb{R}$, $t \in \mathbb{R}^+$. For the choices of α and β , (i.1) describe the nonlinear shallow water waves. Upper and lower bounds are derived for energy decay rate in every finite interval [0,T] which reveals that only the lower bound of the energy decays exponentially.

Key Words: Degasperis-Procesi equation, Camassa-Holm equation, traveling wave

1. INTRODUCTION

The Degasperis-Procesi (D-P) equation

 $u_t - u_{txx} + 4uu_x = 3u_xu_{xx} + uu_{xxx}$ $x \in \mathbb{R}$, t > 0 (1.1) was proposed in [1] as one out of three integral equations within a certain family of third-order nonlinear dispersive partial differential equations; the other two being the well-known Korteweg-de Vries (KdV)

$$u_t - 6uu_x + u_{xxx} = 0 \quad x \in \mathbb{R}, t > 0 \tag{1.2}$$

and Camassa-Holm (C-H) equation [2]

$$u_t - u_{txx} + 3uu_x = 2u_x u_{xx} + uu_{xxx} \quad x \in \mathbb{R}, t > 0$$
(1.3)

which models the shallow water waves.

All weak traveling wave solutions of the D-P equations are classified by Lenells [3]. Similar classification for C-H has also been done in [4]. Degasperis, Holm and Hone [5] investigated D-P equation using the method of asymptotic integrability. This equation has a form similar to C-H shallow water wave equation. The exact integrability of the equation (1.1) investigated in [5]. The solitary wave solutions for modified forms of the equations D-P and C-H are developed by Wazwaz [6].

In this work we are interested in the structural stability of the equations D-P and C-H besides the upper and lower bounds of the energy for these equations. For the structural stability, it is fundamental that one wishes to know if a small change in a coefficient of the equation or boundary data, or small change of the equations themselves will lead to a drastic change in the solution or not. In this article we have proved that

(i.1)

 $u_t - u_{txx} + \alpha uu_x = \beta u_x u_{xx} + uu_{xxx}$ $x \in \mathbb{R}$, t > 0 (1.4) is structurally stable with respect to the coefficients α and β . D-P and C-H equations are attained for the choices $\alpha = 4$, $\beta = 3$ and $\alpha = 3$, $\beta = 2$ respectively. We obtain that upper and lower bounds of the energy for the solutions of equations D-P and C-H are derived in every finite interval [0,T] which shows that only the lower bound of the energy decays exponentially.

2. STRUCTURAL STABILITY

Now we consider the problem

$$u_t - u_{txx} + \alpha u u_x = \beta u_x u_{xx} + u u_{xxx} \quad u \in C_0^{4,1}(\mathbb{R} \times \mathbb{R}^+),$$

$$0 < t < T$$
 for fixed T (2.1)

$$u(x,0) = u_0(x) \qquad x \in \mathbb{R} \tag{2.2}$$

where α , $\beta > 1$ are constants, $C_0^{4,1}(\mathbb{R} \times \mathbb{R}^+)$ is the space of functions having compact support which have fourth order and first order derivative with respect to x and t respectively. To do this, we let (u, α_1, β_1) be the solution of the following problem

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$$u_t - u_{txx} + \alpha_1 u u_x = \beta_1 u_x u_{xx} + u u_{xxx} \quad u \in C_0^{4,1}(\mathbb{R} \times \mathbb{R}^+), \quad 0 < t < T \quad \text{for fixed} \quad T$$

$$(2.3)$$

$$u(x,0) = u_0(x) \qquad x \in \mathbb{R}$$

and (v, α_2, β_2) be the solution of

$$v_t - v_{txx} + \alpha_2 v v_x = \beta_2 v_x v_{xx} + v v_{xxx} \quad v \in C_0^{4,1}(\mathbb{R} \times \mathbb{R}^+), \quad 0 < t < T \quad \text{for fixed} \quad T$$

$$v(x, 0) = u_0(x) \qquad \qquad \mathbb{D}$$

$$(2.5)$$

$$v(x,0) - u_0(x) \qquad x \in \mathbb{R}$$

$$(2.6)$$

where $\alpha_1, \beta_1, \alpha_2, \beta_2 > 1$ are constants. Now, we define the difference of these solutions

by w = u - v, $\alpha = \alpha_1 - \alpha_2$, $\beta = \beta_1 - \beta_2$ where we assume that $\alpha_1 > \alpha_2$ and $\beta_1 > \beta_2$. Then from (2.3)-(2.6), we find that (w, α, β) satisfies the following initial value problem

$$w_t - w_{txx} + \alpha u u_x + \alpha_2 (w u_x + v w_x) - \beta u_x u_{xx} - \beta_2 (w_x u_{xx} + v_x w_{xx}) - (w u_{xxx} + v w_{xxx}) = 0$$
(2.7)

$$w(x,0) = 0 \tag{2.8}$$

We may state our result on structural stability for the problem defined by (2.1)-(2.2) as :

Theorem 1: Let w be the solution of the problem (2.7) and (2.8). Then w satisfies the estimate

$$\|w\|^{2} + 2\|w_{x}\|^{2} + 2\|w_{xx}\|^{2} + \|w_{xxx}\|^{2} \le (\alpha K_{1} + \beta K_{2})(\frac{e^{\gamma I} - 1}{\gamma})$$
(2.9)

for fixed T where K_1 , K_2 and γ are positive constants and $\|\cdot\|$ denotes the L_2 norm of functions. *Proof.* Taking the inner product of (2.7) by w yields

$$(w_t, w) - (w_{txx}, w) + (\alpha u u_x, w) + (\alpha_2 (w u_x + v w_x), w) - (\beta u_x u_{xx}, w) - (\beta_2 (w_x u_{xx} + v_x w_{xx}), w) - (w u_{xxx} + v w_{xxx}, w) = 0$$

which

$$\frac{1}{2}\frac{d}{dt}\left\{\left\|w\right\|^{2}+\left\|w_{x}\right\|^{2}\right\}=-\alpha\int uu_{x}wdx-\alpha_{2}\int (wu_{x}+vw_{x})wdx+\beta\int u_{x}u_{xx}wdx+\beta_{2}\int (w_{x}u_{xx}+v_{x}w_{xx})wdx+\int (wu_{xxx}+vw_{xxx})wdx+(2.11)$$
(2.11)

Since $u, v \in C_0^{4,1}(\mathbb{R} \times \mathbb{R}^+)$ then there exists a generic constant such that the functions u and v together with their derivatives are all bounded by a generic constant D. For the first integral on the right hand side of (2.11) we obtain,

$$-\alpha \int u u_x w dx \le \alpha C \left\{ \left\| w \right\|^2 + \left\| u_x \right\|^2 \right\}$$
(2.12)

utilizing Cauchy and Hölder inequalities. For the second integral in the right hand side of (2.11) we get

$$-\alpha_2 \int (wu_x + vw_x) wdx \le \alpha_2 C \left\{ \|w\|^2 + \|w_x\|^2 \right\}.$$
(2.13)

In a similar way, we can compute the estimates for the other terms as

$$\beta \int u_{x} u_{xx} w dx \le \beta C \left\{ \|w\|^{2} + \|u_{xx}\|^{2} \right\},$$
(2.14)

$$\beta_2 \int (w_x u_{xx} + v_x w_{xx}) w dx \le \beta_2 C \left\{ \|w\|^2 + \|w_x\|^2 + \|w_{xx}\|^2 \right\},$$
(2.15)

$$\int (wu_{xxx} + vw_{xxx})wdx \le C \left\{ \|w\|^2 + \|w_{xxx}\|^2 \right\}.$$
(2.16)

Substituting the estimates (2.12)-(2.16) in (2.11) we find

$$\frac{d}{dt} \left\{ \|w\|^{2} + \|w_{x}\|^{2} \right\} \leq C \left\{ \alpha \|u_{x}\|^{2} + \beta \|u_{xx}\|^{2} \right\} + C \left\{ \alpha + \beta + \alpha_{2} + \beta_{2} + 1 \right\} \|w\|^{2} + C \left\{ (\alpha_{2} + \beta_{2}) \|w_{x}\|^{2} + \beta_{2} \|w_{xx}\|^{2} + \|w_{xxx}\|^{2} \right\}$$

$$(2.17)$$

Similarly taking the inner product of (2.7) by
$$^{w_{XX}}$$
, we find

$$\frac{d}{dt} \left\{ \|w_X\|^2 + \|w_{XX}\|^2 \right\} \leq D \left\{ \alpha \|u_X\|^2 + \beta \|u_{XX}\|^2 \right\} + C \left\{ (\alpha_2 + 1) \|w\|^2 + \alpha_2 \|w_X\|^2 \right\} + \left\{ C (\alpha_2 + \beta_2 + 1) + \beta D \right\} \|w_{XX}\|^2 + C \|w_{XXX}\|^2$$
(2.18)

Now let us differentiate equation of (2.7) with respect to x:

 $w_{tx} - w_{txxx} + \alpha(u_x^2 + uu_{xx}) + \alpha_2(w_xu_x + wu_{xx} + v_xw_x + vw_{xx}) - \beta(u_xu_{xxx} + u_{xx}^2) - \beta_2(w_{xx}u_{xx} + w_xu_{xxx} + v_xw_{xx} + v_xw_{xxx}) - (w_xu_{xxx} + wu_{xxxx} + v_xw_{xxx} + v_xw_{xxx}) = 0$ (2.19)

Taking the inner product of (2.19) by \mathcal{W}_{xxx} , we find

$$\frac{d}{dt} \left\{ \|w_{xx}\|^{2} + \|w_{xxx}\|^{2} \right\} \leq C \left\{ \alpha \|u_{xx}\|^{2} + \beta \|u_{xxx}\|^{2} \right\} + C \left\{ (\alpha_{2} + 1) \|w\|^{2} + (\alpha_{2} + \beta_{2} + 1) \|w_{x}\|^{2} + (\alpha + \alpha_{2} + \beta + \beta_{2}) \|w_{xx}\|^{2} \right\} + C \left\{ (\alpha + \alpha_{2} + \beta + \beta_{2} + 1) \|w_{xxx}\|^{2} \right\}.$$

$$(2.20)$$

Adding up the inequalities (2.17), (2.18), (2.20) we have

$$\frac{d}{dt} \left\{ \|w\|^{2} + 2\|w_{xx}\|^{2} + 2\|w_{xxx}\|^{2} + \|w_{xxx}\|^{2} \right\} \le \alpha K_{1} + \beta K_{2} + \gamma \left\{ \|w\|^{2} + 2\|w_{xx}\|^{2} + 2\|w_{xxx}\|^{2} + \|w_{xxx}\|^{2} \right\}$$
(2.21)

where
$$K_1 = (C+D) \|u_x\|^2 + C \|u_{xx}\|^2$$
, $K_2 = (C+D) \|u_{xx}\|^2 + C \|u_{xxx}\|^2$ and
 $\gamma = C \max\{ (\alpha + \beta + 3\alpha_2 + 3 + \beta_2), (3\alpha_2 + 1 + 2\beta_2), (\alpha + \beta + \alpha_2 + 2\beta_2), (\alpha + \beta + \alpha_2 + \beta_2 + 2) \}$
Thus, from (2.20), we have

Thus, from (2.20) we have

$$\frac{d}{dt}\Psi(t) - \gamma \Psi(t) \le \alpha K_1 + \beta K_2 \tag{2.22}$$

where $\Psi(t) = \|w\|^2 + 2\|w_x\|^2 + 2\|w_{xx}\|^2 + \|w_{xxx}\|^2$. Solving the differential inequality (2.22), we arrive at

$$\|w\|^{2} + 2\|w_{x}\|^{2} + 2\|w_{xx}\|^{2} + \|w_{xxx}\|^{2} \le (\alpha K_{1} + \beta K_{2}) \left(\frac{e^{\gamma T} - 1}{\gamma}\right)$$

for fixed T. And so we have completed the proof of the theorem.

Remark 2. W and its X derivatives of order up to ³, tends to zero as $\alpha \to 0$, $\beta \to 0$ for finite T

So, the solutions of (2.22) depend continuously on α and β in $C_0^{4,1}(\mathbb{R} \times \mathbb{R}^+)$ which means that (2.1)-(2.2) is structurally stable with respect to the coefficients α and β .

3. UPPER AND LOWER BOUNDS ON THE ENERGY

Let u be a solution to the initial-value problem (2.1), (2.2) with $\alpha > 1$ and $\beta > 1$. By similar computations given in Section 2, we find

$$\frac{d}{dt} \left\{ \left\| u \right\|^2 + 2 \left\| u_x \right\|^2 + 2 \left\| u_{xxx} \right\|^2 + \left\| u_{xxx} \right\|^2 \right\} = (2 - \beta - \alpha) \int u^2 u_{xxx} dx - (2\beta + 1) \int u_x u_{xxx}^2 dx + (1 - 2\beta - 5\alpha) \int u_x u_{xx}^2 dx$$
(3.1)

If we use Cauchy and Hölder inequalities in (3.1) we obtain

$$\frac{d}{dt} \left\{ \left\| u \right\|^{2} + 2 \left\| u_{x} \right\|^{2} + 2 \left\| u_{xx} \right\|^{2} + \left\| u_{xxx} \right\|^{2} \right\} \ge (1 - \frac{\beta}{2} - \frac{\alpha}{2}) \max u^{2} \left\| u \right\|^{2} + 2(-\frac{5\alpha}{4} - \beta) \left\| u_{x} \right\|^{2} + 2\left(\frac{1}{4} - \frac{\beta}{2} - \frac{5\alpha}{4}\right) \max u^{2}_{xx} \left\| u_{xx} \right\|^{2} + \left(1 - \frac{\beta}{2} - \frac{\alpha}{2} - \left(\beta + \frac{1}{2}\right) \max u^{2}_{xxx} \right) \left\| u_{xxx} \right\|^{2}$$

$$(3.2)$$

Taking

$$\eta = \max\left\{ \left(1 - \frac{\beta}{2} - \frac{\alpha}{2}\right) \max u^2, \left(-\frac{5\alpha}{4} - \beta\right), \left(\frac{1}{4} - \frac{\beta}{2} - \frac{5\alpha}{4}\right) \max u^2_{xx}, \left(1 - \frac{\beta}{2} - \frac{\alpha}{2} - \left(\beta + \frac{1}{2}\right) \max u^2_{xxx}\right) \right\}$$

and

$$Y(t) = \|u\|^2 + 2\|u_x\|^2 + 2\|u_{xx}\|^2 + \|u_{xxx}\|^2,$$

we have

$$\frac{d}{dt}Y(t) - \eta Y(t) \ge 0$$
(3.3)

Solving the inequality (3.3) we arrive at

$$e^{\eta T} \left\{ \left\| u(x,0) \right\|^{2} + 2 \left\| u_{x}(x,0) \right\|^{2} + 2 \left\| u_{xx}(x,0) \right\|^{2} + \left\| u_{xxx}(x,0) \right\|^{2} \right\} \le \left\| u(x,t) \right\|^{2} + 2 \left\| u_{x}(x,t) \right\|^{2} + 2 \left\| u_{xxx}(x,t) \right\|^{2} + \left\| u_{xxx}(x,t) \right\|^{2}$$
(3.4)

where $\eta \leq 0$. This inequality gives a lower bound for the energy.

Now we will derive an upper bound for the energy. From (3.1) we have

$$\frac{d}{dt} \left\{ \|u\|^{2} + 2\|u_{xx}\|^{2} + 2\|u_{xx}\|^{2} + \|u_{xxx}\|^{2} \right\} \leq \left| 1 - \frac{\beta + \alpha}{2} \right| \max |u| \|u\|^{2} + 2 \left\{ \left| \frac{1 - 5\alpha - 2\beta}{4} \right| \max |u_{xx}| + \left(\frac{\beta}{2} + \frac{1}{4} \right) \max |u_{xxx}| \right\} \|u_{xx}\|^{2} + 2 \left| \frac{1 - 5\alpha - 2\beta}{4} \right| \max |u_{xx}| + \left| \frac{\beta + \alpha}{2} \right| \max |u_{xxx}|^{2} + 2 \left| \frac{1 - 5\alpha - 2\beta}{4} \right| \max |u_{xxx}|^{2} + \left\{ \left| 1 - \frac{\beta + \alpha}{2} \right| \max |u| + \left(\beta + \frac{1}{2} \right) \max |u_{xxx}| \right\} \|u_{xxx}\|^{2} \right\}$$
(3.5)

Taking

$$\mu = \max\left\{ \left| 1 - \frac{\beta + \alpha}{2} \right| \max |u|, \left| \frac{1 - 5\alpha - 2\beta}{4} \right| \max |u_{xx}| + \left(\frac{\beta}{2} + \frac{1}{4}\right) \max |u_{xxx}|, \left| \frac{1 - 5\alpha - 2\beta}{4} \right| \max |u_{xx}|, \left| \frac{1 - \beta - 2\beta}{4} \right| \max |u_{xx}|, \left| \frac{1 - \beta - 2\beta}{4} \right| \max |u_{xxx}| \right\} \right\}$$

and

$$Y(t) = \|u\|^2 + 2\|u_x\|^2 + 2\|u_{xx}\|^2 + \|u_{xxx}\|^2,$$

we have

$$\frac{d}{dt}Y(t) - \mu Y(t) \le 0 \tag{3.6}$$

Then integrating the inequality (3.6) from 0 to *T* we arrive at
$$\|u(x,T)\|^{2} + 2\|u_{x}(x,T)\|^{2} + 2\|u_{xx}(x,T)\|^{2} + \|u_{xxx}(x,T)\|^{2} \le e^{\mu T} \left\{ \|u(x,0)\|^{2} + 2\|u_{x}(x,0)\|^{2} + 2\|u_{xx}(x,0)\|^{2} + \|u_{xxx}(x,0)\|^{2} \right\}$$
(3.7)

where $\mu \ge 0$. (3.7) gives an upper bound for the energy in every finite interval [0,T].

We may combine the above results as in the following theorem.

Theorem 3. The energy corresponding to the solutions of the initial value problem (2.1)-(2.2) in $C_0^{4,1}(\mathbb{R} \times \mathbb{R}^+)$ satisfy

$$e^{\eta T} \left\{ \left\| u(x,0) \right\|^{2} + 2 \left\| u_{x}(x,0) \right\|^{2} + 2 \left\| u_{xx}(x,0) \right\|^{2} + \left\| u_{xxx}(x,0) \right\|^{2} \right\} \le \left\| u(x,T) \right\|^{2} + 2 \left\| u_{x}(x,T) \right\|^{2} + 2 \left\| u_{xx}(x,T) \right\|^{2} + \left\| u_{xxx}(x,T) \right\|^{2} \\ \le e^{\mu T} \left\{ \left\| u(x,0) \right\|^{2} + 2 \left\| u_{x}(x,0) \right\|^{2} + 2 \left\| u_{xx}(x,0) \right\|^{2} + \left\| u_{xxx}(x,0) \right\|^{2} \right\}$$

For fixed T where α , $\beta > 1$ are constants.

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