



Ambarzumyan-Type Theorem for a Conformable Fractional Diffusion Operator

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Abstract

In this paper, we prove an Ambarzumyan-type theorem for a Conformable fractional diffusion operator, i.e. we show that q(x) and p(x) functions are zero if the eigenvalues are the same as the eigenvalues of zero potentials.

Keywords: Ambarzumyan-type theorem, Conformable fractional derivative, Diffusion operator, Inverse problem **2010 AMS:** 34A55, 34B24

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1. Introduction

Inverse problems of spectral analysis consist of recovering the coefficients of an operator from their spectral characteristics. Such problems often appear in mathematics, mathematical physics, mechanics, and other branches of natural sciences.

The first result in inverse spectral problems for the Sturm-Liouville operator has been obtained by Ambarzumyan in 1929 (see [1]). He considered the boundary value problem

$$\begin{cases} -y'' + q(x)y = \lambda y, & 0 < x < \pi \\ y'(0) = y'(\pi) = 0 \end{cases}$$
(1.1)

and proved that if the eigenvalues of (1.1) are $\lambda_n = n^2$, $n \ge 0$, then $q(x) \equiv 0$ a.e. on $(0, \pi)$.

Until now, some Ambarzumyan-type theorems for the Sturm-Liouville, Dirac and diffusion operators including classical derivatives have been studied by many authors (see [2]-[17]). Particularly, in [15], by considering a quadratic Sturm-Liouville operator called the diffusion operator, it is shown that q(x) and p(x) functions are zero if the spectrum is the same as the spectrum of zero potential.

In 2014, Khalil et al. gave a new fractional derivative definition called conformable derivative that extends the well-known limit definition of the classical derivative (see [18]). The conformable fractional derivative has some advantages over fractional derivatives. For instance, while some properties such as the derivative of the product of two functions, the derivative of the quotient of two functions, and the chain rule are not satisfied in all fractional derivatives, these properties are satisfied in the conformable fractional derivative were developed in the works in [19]-[24].

In recent years, the direct and inverse problems for the various operators which include conformable fractional derivative have been studied (see [25]-[33]). These problems appear in various branches of applied sciences (see [34]-[38]). In the current

literature, there are no results related to the Ambarzumyan-type theorem for a diffusion operator with the conformable fractional derivative.

We consider a conformable fractional diffusion operator (CFDO) with Neumann boundary conditions. The operator $L_{\alpha} = L_{\alpha}(p(x), q(x))$ is the form

$$\begin{cases} -T_x^{\alpha} T_x^{\alpha} y + [2\lambda p(x) + q(x)] y = \lambda^2 y, \quad 0 < x < \pi \\ T_x^{\alpha} y(0) = T_x^{\alpha} y(\pi) = 0 \end{cases}$$
(1.2)

where λ is the spectral parameter, $\alpha \in (0,1]$, $q(x) \in W_{2,\alpha}^1[0,\pi]$, $p(x) \in W_{2,\alpha}^2[0,\pi]$ are real-valued functions and $T_x^{\alpha} y$ is a conformable fractional derivative of order α of y at x.

The goal of this paper is to prove an Ambarzumyan-type theorem for the operator L_{α} . The result obtained can be considered as a partial α -generalization of the result given in [15].

2. Preliminaries

Definition 2.1. Let $f : [0,\infty) \to \mathbb{R}$ be a given function. Then, the conformable fractional derivative of f of order α with respect to x is defined by

$$T_{x}^{\alpha}f(x) = \lim_{h \to 0} \frac{f(x + hx^{1 - \alpha}) - f(x)}{h}, \ T_{x}^{\alpha}f(0) = \lim_{x \to 0^{+}} T_{x}^{\alpha}f(x)$$

for all x > 0, $\alpha \in (0,1]$. If f is differentiable that is $f'(x) = \lim_{h \to 0} \frac{f(x+h) - f(x)}{h}$, then, $T_x^{\alpha} f(x) = x^{1-\alpha} f'(x)$.

We note that more detailed knowledge about conformable fractional calculus can be seen in [18]-[24]. Let $\varphi(x, \lambda; \alpha)$ be the solution of equation (1.2) satisfying the initial conditions

$$\varphi(0,\lambda;\alpha) = 1, T_x^{\alpha}\varphi(0,\lambda;\alpha) = 0$$

From [32], this solution can be shown with the α -integral representation for h = H = 0 as

$$\varphi(x,\lambda;\alpha) = \cos\left(\lambda \frac{x^{\alpha}}{\alpha} - \theta(x)\right) + \int_{0}^{x} A\left(x, \frac{t^{\alpha}}{\alpha}\right) \cos\lambda \frac{t^{\alpha}}{\alpha} d_{\alpha}t + \int_{0}^{x} B\left(x, \frac{t^{\alpha}}{\alpha}\right) \sin\lambda \frac{t^{\alpha}}{\alpha} d_{\alpha}t,$$
(2.1)

where the functions $A\left(x, \frac{t^{\alpha}}{\alpha}\right)$ and $B\left(x, \frac{t^{\alpha}}{\alpha}\right)$ satisfy the following system

$$\begin{cases} T_{x}^{\alpha}T_{x}^{\alpha}A\left(x,\frac{t^{\alpha}}{\alpha}\right) - q\left(x\right)A\left(x,\frac{t^{\alpha}}{\alpha}\right) - 2p\left(x\right)T_{t}^{\alpha}B\left(x,\frac{t^{\alpha}}{\alpha}\right) = T_{t}^{\alpha}T_{t}^{\alpha}A\left(x,\frac{t^{\alpha}}{\alpha}\right) \\ T_{x}^{\alpha}T_{x}^{\alpha}B\left(x,\frac{t^{\alpha}}{\alpha}\right) - q\left(x\right)B\left(x,\frac{t^{\alpha}}{\alpha}\right) + 2p\left(x\right)T_{t}^{\alpha}A\left(x,\frac{t^{\alpha}}{\alpha}\right) = T_{t}^{\alpha}T_{t}^{\alpha}B\left(x,\frac{t^{\alpha}}{\alpha}\right). \end{cases}$$

Besides, the following relations are provided:

$$B(x,0) = 0, \quad T_t^{\alpha} A\left(x, \frac{t^{\alpha}}{\alpha}\right)\Big|_{t=0} = 0,$$

$$\theta(x) = \int_0^x p(t) d_{\alpha} t,$$

$$A\left(0,0\right) = 0,$$

$$A\left(x, \frac{x^{\alpha}}{\alpha}\right) \cos \theta(x) + B\left(x, \frac{x^{\alpha}}{\alpha}\right) \sin \theta(x) = \frac{1}{2} \int_0^x \left(q(t) + p^2(t)\right) d_{\alpha} t,$$

$$\theta(x) = p(0) \frac{x^{\alpha}}{\alpha} + 2 \int_0^x \left[A\left(s, \frac{s^{\alpha}}{\alpha}\right) \sin \theta(s) - B\left(s, \frac{s^{\alpha}}{\alpha}\right) \cos \theta(s)\right] d_{\alpha} s.$$

The function

$$\Delta_{\alpha}(\lambda) = T_{x}^{\alpha} \varphi(\pi, \lambda; \alpha)$$
(2.2)

is entire function in λ and is called as the characteristic function of operator L_{α} . It is well-known that the roots of $\Delta_{\alpha}(\lambda) = 0$ are coincide with the eigenvalues of operator L_{α} .

From (2.1) and (2.2), we have easily that

$$\Delta_{\alpha}(\lambda) = -(\lambda - p(\pi))\sin\left(\lambda\frac{\pi^{\alpha}}{\alpha} - \theta(\pi)\right) + A\left(\pi, \frac{\pi^{\alpha}}{\alpha}\right)\cos\lambda\frac{\pi^{\alpha}}{\alpha} + B\left(\pi, \frac{\pi^{\alpha}}{\alpha}\right)\sin\lambda\frac{\pi^{\alpha}}{\alpha} + \int_{0}^{\pi} \left(T_{x}^{\alpha}A\left(\pi, \frac{t^{\alpha}}{\alpha}\right)\right)\cos\lambda\frac{t^{\alpha}}{\alpha}d_{\alpha}t + \int_{0}^{\pi} \left(T_{x}^{\alpha}B\left(\pi, \frac{t^{\alpha}}{\alpha}\right)\right)\sin\lambda\frac{t^{\alpha}}{\alpha}d_{\alpha}t.$$
(2.3)

With the help of (2.3), the following theorem for the eigenvalues λ_n of operator L_{α} can be proved as in [32]:

Theorem 2.2. The operator L_{α} has a countable set of eigenvalues $\{\lambda_n\}$ and the following asymptotic formula holds:

$$\lambda_n = \frac{n\alpha}{\pi^{\alpha-1}} + c_{\alpha,0} + \frac{c_{\alpha,1}}{n} + o\left(\frac{1}{n}\right), \ |n| \to \infty,$$
(2.4)

where

$$c_{\alpha,0} = \frac{\alpha}{\pi^{\alpha}} \int_{0}^{\pi} p(x) d_{\alpha}x, \ c_{\alpha,1} = \frac{1}{2\pi} \int_{0}^{\pi} \left(q(x) + p^{2}(x)\right) d_{\alpha}x.$$

3. Main Result

In this section, we prove an Ambarzumyan-type theorem for the operator L_{α} , i.e. we show that q(x) and p(x) functions are zero if the eigenvalues are the same as the eigenvalues of zero potentials.

Theorem 3.1. If the eigenvalues of the operator L_{α} are $\lambda_n = \frac{n\alpha}{\pi^{\alpha-1}}$, $n \in \mathbb{Z}$, then for each fixed α , q(x) = 0, p(x) = 0 a.e. on $(0,\pi)$ and $\theta(\pi) = 0$.

Proof. It follows from (2.4) that for each fixed α , $c_{\alpha,0} = 0$, $c_{\alpha,1} = 0$, i.e.

$$\int_{0}^{\pi} p(x) d_{\alpha} x = 0 = \theta(\pi)$$

and

$$\int_{0}^{\pi} q(x) d_{\alpha} x = -\int_{0}^{\pi} p^{2}(x) d_{\alpha} x.$$
(3.1)

Let $y_0(x; \alpha) = y(x, 0; \alpha)$ be an eigenfunction corresponding to the eigenvalue $\lambda_0 = 0$ of the operator L_{α} . Then we can write

$$\begin{cases} -T_x^{\alpha} T_x^{\alpha} y_0(x; \alpha) + q(x) y_0(x; \alpha) = 0, & 0 < x < \pi \\ T_x^{\alpha} y_0(0; \alpha) = 0, & T_x^{\alpha} y_0(\pi; \alpha) = 0 \end{cases}$$
(3.2)

It is clear that $y_0(0; \alpha) \neq 0$ and $y_0(\pi; \alpha) \neq 0$. Otherwise, $y_0(0; \alpha) = T_x^{\alpha} y_0(0; \alpha) = 0$ or $y_0(\pi; \alpha) = T_x^{\alpha} y_0(\pi; \alpha) = 0$. In both cases, we get $y_0(x; \alpha) = 0$ through the uniqueness of the solution of an initial value problem, which contradicts the fact that $y_0(x; \alpha)$ is an eigenfunction.

Taking into account the relation

$$\frac{T_x^{\alpha}T_x^{\alpha}y_0(x;\alpha)}{y_0(x;\alpha)} = T_x^{\alpha}\left(\frac{T_x^{\alpha}y_0(x;\alpha)}{y_0(x;\alpha)}\right) + \left(\frac{T_x^{\alpha}y_0(x;\alpha)}{y_0(x;\alpha)}\right)^2,$$

we obtain from (3.2) that

$$T_{x}^{\alpha}\left(\frac{T_{x}^{\alpha}y_{0}(x;\alpha)}{y_{0}(x;\alpha)}\right)+\left(\frac{T_{x}^{\alpha}y_{0}(x;\alpha)}{y_{0}(x;\alpha)}\right)^{2}=q(x).$$

By α -integrating of both sides of the above equation from 0 to π

$$\int_{0}^{\pi} T_{x}^{\alpha} \left(\frac{T_{x}^{\alpha} y_{0}\left(x;\alpha\right)}{y_{0}\left(x;\alpha\right)} \right) d_{\alpha}x + \int_{0}^{\pi} \left(\frac{T_{x}^{\alpha} y_{0}\left(x;\alpha\right)}{y_{0}\left(x;\alpha\right)} \right)^{2} d_{\alpha}x = \int_{0}^{\pi} q(x) d_{\alpha}x$$

is obtained. From (3.1) and (3.2), we get

$$\frac{T_x^{\alpha} y_0(x;\alpha)}{y_0(x;\alpha)} \Big|_0^{\pi} + \int_0^{\pi} \left(\frac{T_x^{\alpha} y_0(x;\alpha)}{y_0(x;\alpha)}\right)^2 d_{\alpha} x = -\int_0^{\pi} p^2(x) d_{\alpha} x$$

or

$$\int_{0}^{\pi} \left[\left(\frac{T_{x}^{\alpha} y_{0}\left(x;\alpha\right)}{y_{0}\left(x;\alpha\right)} \right)^{2} + p^{2}\left(x\right) \right] d_{\alpha}x = 0.$$

Thus, for each fixed α , $(T_x^{\alpha} y_0(x; \alpha))^2 + p^2(x) y_0^2(x; \alpha) \equiv 0$, i.e. p(x) = 0 a.e. on $(0, \pi)$ and $y_0(x; \alpha) \equiv k$, where $0 \neq k - const$. Hence, it is concluded from (3.2) that

$$-T_x^{\alpha}T_x^{\alpha}k + q(x)k = 0$$

then q(x) = 0 a.e. on $(0, \pi)$. Therefore, the theorem is completed.

4. Conclusion

As known, the inverse problems of spectral analysis consist in recovering operators from their spectral characteristics, and the first result in this direction belongs to Ambarzumyan for the Sturm-Liouville operator. Until today, many studies have been done on the Ambarzumyan theorem for various operators including classical derivatives. In this paper, the Ambarzumyan theorem is proved for the diffusion operator with Neumann boundary value conditions including fractional derivatives. This study will make an important contribution to the inverse problems of spectral analysis. This theorem can be proved in the future for various operators with different derivatives.

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