Konuralp Journal of Mathematics, 11 (2) (2023) 195-205



Konuralp Journal of Mathematics

Research Paper Journal Homepage: www.dergipark.gov.tr/konuralpjournalmath e-ISSN: 2147-625X



(2.1)

(2.2)

Finding Powerful Solutions for the Generalized Hyperelastic-Rod Wave Equation

Şeyma Tülüce Demiray¹ and Uğur Bayrakçı^{1*}

¹Department of Mathematics, Faculty of Science and Letters, Osmaniye Korkut Ata University, 80000, Osmaniye, Türkiye *Corresponding author

Abstract

In this paper, the generalized hyperelastic rod wave equation has been studied. The generalized exponential rational function method (GERFM) has been applied to the generalized hyperelastic rod wave equation. Thus, some new and abundant soliton solutions of the generalized hyperelastic rod wave equation have been obtained. Also, in Wolfram Mathematica 12, both 2D and 3D shapes of these built-in results have been plotted.

Keywords: GERFM, *Generalized hyperelastic-rod wave equation, soliton solutions.* 2010 Mathematics Subject Classification: 35C07, 35A20, 35A25.

1. Introduction

In this study, GERFM have been used, the solution methods of nonlinear evolution equations (NLEEs), and these methods have been applied to the generalized hyperelastic-rod wave equation, which is a variant of NLEEs. NLEEs have very important applications in areas such as mathematical physics, optical fibres, mathematical chemistry, hydrodynamics, fluid dynamics, geochemistry, control theory, meteorology, optics, mechanics, chemical kinematics, biophysics, biogenetics, and so on. A number of methods have been developed by various researchers in order to obtain solutions of NLEEs, which have such important areas of use in the scientific World [1, 2, 3, 4, 5, 6, 7, 8, 9, 10, 11, 12, 13, 14, 15]. Generalized hyperelastic-rod wave equation is given as [16, 17]:

$$u_t - u_{xxt} - \alpha u_x + 2\beta u u_x + 3\theta u^2 u_x - \gamma u_x u_{xx} - u u_{xxx} = 0,$$

$$\tag{1.1}$$

where α, β, θ and γ are constants and we accept that θ is nonzero. This equation is used to describe finite length, small amplitude radial deformation waves in cylindrical compressible hyperelastic rods. This equation also includes important physical models. For $\beta = \frac{3}{2}, \theta = 0$ and $\gamma = 2$, the equation can be reduced to Camassa-Holm equation. For $\alpha = 1, \beta = \frac{1}{2}, \theta = 0$ and $\gamma = 3$, the equation can be reduced to Fornberg-Whitham equation. Also, for $\beta = 2, \theta = 0$ and $\gamma = 3$, the equation can be reduced to Degasperis-Procesi equation [16, 17, 18, 19, 20].

Generalized hyperelastic-rod wave equation has been studied by some researchers recently. Akcagil et al. got the travelling wave solutions of the equation with the help of expansion method [16]. Gözükızıl and Akçagıl obtained the exact solutions of the equation using the tanh-coth method [17].

This study, which was prepared to specified the soliton solutions of the generalized hyperelastic-rod wave equation using GERFM [21, 22, 23, 24, 25], was designed as follows: In Section 2, GERFM's basic principles are presented. In Section 3, some soliton solutions of generalized hyperelastic-rod wave equation have been obtained by applying methods.

2. Definition of GERFM

Step1: We consider NLPDE given below:

 $P(u, u_x, u_t, u_{xx}, \ldots) = 0.$

We first apply the wave transform given as below to Eq. (2.1);

 $u(x,t) = u(\eta), \eta = k(x-\lambda)t.$

Email addresses: ubayrakci42@gmail.com (U. Bayrakçı), seymatuluce@gmail.com (§ Tülüce Demiray)

(3.4)

Eq. (2.1) is transformed into ordinary differential equation by using Eq. (2.2):

$$R(u, u', u'', \cdots) = 0, (2.3)$$

where *k* and λ values that are not taken into account will be calculated later. Step2: Assume that we think the solutions of Eq. (2.3) as:

$$u(\eta) = a_0 + \sum_{i=1}^{M} a_i \Phi(\eta)^i + \sum_{i=1}^{M} \frac{b_i}{\Phi(\eta)^i},$$
(2.4)

where

$$\Phi(\eta) = \frac{p_1 e^{q_1 \eta} + p_2 e^{q_2 \eta}}{p_3 e^{q_3 \eta} + p_4 e^{q_4 \eta}}.$$
(2.5)

Here value of *M* is determined through the homogeneous balance principle. $p_1, p_2, p_3, p_4, q_1, q_2, q_3, q_4, a_1, a_2, \dots a_M, b_1, b_2, \dots b_M$ constants are determined to fit the solution.

Step3: If the Eq. (2.4) is taken into account in the Eq. (2.3), $P(e^{q_1\eta}, e^{q_2\eta}, e^{q_3\eta}, e^{q_4\eta}) = 0$ equation system is obtained. A system of equations is obtained by equating all coefficients of *P* to zero.

Step4: If we solve the obtained system of equations and the found values consider in Eq. (2.4), the solutions of the discussed NLPDE are obtained.

3. Application of GERFM

To find the exact solutions of Eq. (1.1) we consider the following transformation:

$$u(x,t) = u(\eta), \eta = x - vt.$$

$$(3.1)$$

Replace Eq. (3.1) into Eq. (1.1) and we get,

$$vu' + vu''' + \alpha u' + 2\beta uu' + 3\theta u^2 u' - \gamma u'u'' - uu''' = 0.$$
(3.2)

Integrating Eq. (3.2) and if the integration constant is set to zero, we have

$$(\alpha - v)u + vu'' + \beta u^2 + \theta u^3 - \frac{\gamma - 1}{2}(u')^2 - uu'' = 0.$$
(3.3)

By using balance principle in Eq. (3.3), we obtain

$$M = 2.$$

If M = 2 is taken into account in Eq. (2.4).

$$\mu(\eta) = a_0 + a_1 \Phi(\eta) + a_2 \Phi^2(\eta) + \frac{b_1}{\Phi(\eta)} + \frac{b_2}{\Phi^2(\eta)},$$
(3.5)

equality is achieved. So the obtained different states of the considered equation via GERFM are as follows: Family 1: For p = [-1, 0, 1, 1] and q = [0, 0, 1, 0], Eq. (2.5) turns into the form,

$$\Phi(\eta) = \frac{-1}{1+e^{\eta}}.\tag{3.6}$$

Case1:

$$a_{0} = \frac{-1 + 2\beta + A}{6\theta}, a_{1} = \frac{6 - 3A + \sqrt{3}\sqrt{7 + 12\alpha\theta} - 4A + 4\beta(2 + \beta - A)}{3\theta}, b_{1} = 0, b_{2} = 0,$$

$$a_{2} = \frac{6 - 3A + \sqrt{3}\sqrt{7 + 12\alpha\theta} - 4A + 4\beta(2 + \beta - A)}{3\theta}, v = \frac{-1 - 12\alpha\theta + A + \beta(1 + 2\beta + A)}{18\theta},$$

$$\gamma = \frac{1}{6} \left(-6 - 3A + \sqrt{3}\sqrt{7 + 12\alpha\theta} - 4A + 4\beta(2 + \beta - A)}\right)$$
(3.7)

where $A = \sqrt{(1-2\beta)^2 - 12\alpha\theta}$. Embedding the these values in Eqs. (3.5) and (3.6), we acquire the soliton solution of Eq. (1.1)

$$u_{1}(x,t) = -\frac{e^{(x-vt)} \left(5 + 2\beta - 2A\sqrt{3}\sqrt{7 + 12\alpha\theta - 4A + 4\beta(2+\beta-A)}\right) + (-1 + 2\beta + A)\cosh[x-vt]}{3\left(1 + e^{(x-vt)}\right)^{2}\theta},$$
(3.8)

where $v = \frac{-1 - 12\alpha\theta + A + \beta(1 + 2\beta + A)}{18\theta}$.



Figure 3.1: 3D plot of solution (3.8) for $\alpha = 2, \beta = 0.3, \theta = 0.5$ values with $-25 \le x \le 25, -2 \le t \le 2$ range and 2D plot of solution for t = 1.5 with these values.

Case2:

$$a_{0} = -\frac{2\alpha}{-2+\beta}, a_{1} = -\frac{8\alpha(-1+\beta)}{(-2+\beta)\beta}, b_{1} = 0, b_{2} = 0,$$

$$a_{2} = -\frac{8\alpha(-1+\beta)}{(-2+\beta)\beta}, \theta = \frac{4\alpha}{(-2+\beta)\beta}, v = -\frac{2\alpha}{-2+\beta}, \gamma = -1-\beta.$$
(3.9)

Embedding the these values in Eqs. (3.5) and (3.6), we acquire the soliton solution of Eq. (1.1)

$$u_{2}(x,t) = -\frac{\alpha \left(-2 + \beta - \beta \cosh\left[x + \frac{2\alpha t}{-2 + \beta}\right]\right) \operatorname{sech}^{2}\left[\frac{x}{2} + \frac{\alpha t}{-2 + \beta}\right]}{(-2 + \beta)\beta}.$$
(3.10)



Figure 3.2: 3D plot of solution (3.10) for $\alpha = 1.5$, $\beta = 3$ values with $-20 \le x \le 20$, $-4 \le t \le 4$ range and 2D plot of solution for t = 3 with these values.

Case3:

$$a_{0} = \frac{3}{5\theta}, a_{1} = \frac{1}{5\theta}, a_{2} = \frac{1}{5\theta}, b_{1} = 0, b_{2} = 0,$$

$$\alpha = \frac{753}{850\theta}, \beta = -\frac{387}{340}, \theta = \frac{4\alpha}{(-2+\beta)\beta}, v = \frac{957}{1700\theta}, \gamma = -\frac{19}{10}.$$
 (3.11)

Embedding the these values in Eqs. (3.5) and (3.6), we acquire the soliton solution of Eq. (1.1)

$$u_{3}(x,t) = \frac{5 + 6\cosh\left[x - \frac{957t}{1700\theta}\right]}{10\theta + 10\theta\cosh\left[x - \frac{957t}{1700\theta}\right]}.$$
(3.12)



Figure 3.3: 3D plot of solution (3.12) for $\theta = 1.5$ values with $-10 \le x \le 10, -5 \le t \le 5$ range and 2D plot of solution for t = 2 with these values.

Case4:

$$a_{0} = \frac{1}{3\theta}, a_{1} = -\frac{1}{3\theta}, a_{2} = -\frac{1}{3\theta}, b_{1} = 0, b_{2} = 0,$$

$$\alpha = \frac{67}{126\theta}, \beta = -\frac{67}{84}, v = \frac{95}{252\theta}, \gamma = -\frac{13}{6}.$$
(3.13)

. -

Embedding the these values in Eqs. (3.5) and (3.6), we acquire the soliton solution of Eq. (1.1)

$$u_4(x,t) = \frac{4 + \operatorname{sech}^2 \left[\frac{1}{2} \left(x - \frac{95t}{252\theta} \right) \right]}{12\theta}.$$
(3.14)



Figure 3.4: 3D plot of solution (3.14) for $\theta = 2$ values with $-20 \le x \le 20, -1 \le t \le 1$ range and 2D plot of solution for t = 0.5 with these values.

Case5:

$$a_{0} = \frac{3 - i\sqrt{3}}{6\theta}, a_{1} = \frac{-i}{\sqrt{3}\theta}, a_{2} = \frac{-i}{\sqrt{3}\theta}, b_{1} = 0, b_{2} = 0,$$

$$\alpha = \frac{271 - 81i\sqrt{3}}{372\theta}, \beta = \frac{1}{372} \left(-357 + 86i\sqrt{3}\right), v = \frac{395 - 81i\sqrt{3}}{744\theta}, \gamma = -2 - \frac{i}{2\sqrt{3}}.$$
 (3.15)

Embedding the these values in Eqs. (3.5) and (3.6), we acquire the soliton solution of Eq. (1.1)

$$u_{5}(x,t) = \frac{3 + 3\cosh\left[x - \frac{(395 - 81i\sqrt{3})t}{744\theta}\right] - i\sqrt{3}\cosh\left[x - \frac{(395 - 81i\sqrt{3})t}{744\theta}\right]}{6\theta + 6\theta\cosh\left[x - \frac{(395 - 81i\sqrt{3})t}{744\theta}\right]}.$$
(3.16)



Figure 3.5: 3D plot of solution (3.16) for $\theta = 0.5$ values with $-15 \le x \le 15, -4 \le t \le 4$ range and 2D plot of solution for t = 3 with these values.

Family 2: For p = [-2 - i, 2 - i, -1, 1] and q = [i, -i, i, -i], Eq. (2.5) turns into the form,

$$\Phi(\eta) = \frac{\cos(\eta) + 2\sin(\eta)}{\sin(\eta)}.$$
(3.17)

Case1:

$$a_{0} = \frac{10(2+\gamma)}{\theta}, a_{1} = 0, a_{2} = 0, b_{1} = -\frac{40(2+\gamma)}{\theta}, b_{2} = -\frac{50(2+\gamma)}{\theta}, \\ \alpha = -\frac{5(2+\gamma)(2+\beta+2\gamma)}{3\theta}, v = -\frac{(2+\gamma)(2+\beta+2\gamma)}{3\theta}.$$
(3.18)

Embedding the these values in Eqs. (3.5) and (3.17), we acquire the soliton solution of Eq. (1.1)

$$u_6(x,t) = \frac{10(2+\gamma)}{\theta\left(\cos\left[x + \frac{t(2+\gamma)(2+\beta+2\gamma)}{3\theta}\right] + 2\sin\left[x + \frac{t(2+\gamma)(2+\beta+2\gamma)}{3\theta}\right]\right)^2}.$$
(3.19)



Figure 3.6: 3D plot of solution (3.19) for $\gamma = 2, \beta = 3, \theta = 0.5$ values with $-20 \le x \le 20, -1 \le t \le 1$ range and 2D plot of solution for t = 0.5 with these values.

Case2:

$$a_{0} = \frac{10(2+\gamma)}{\theta}, a_{1} = -\frac{8(2+\gamma)}{\theta}, a_{2} = -\frac{2(2+\gamma)}{\theta}, b_{1} = 0, b_{2} = 0,$$

$$\alpha = -\frac{5(2+\gamma)(2+\beta+2\gamma)}{3\theta}, \nu = -\frac{(2+\gamma)(2+\beta+2\gamma)}{3\theta}.$$
 (3.20)

Embedding the these values in Eqs. (3.5) and (3.17), we acquire the soliton solution of Eq. (1.1)

$$u_7(x,t) = \frac{2(2+\gamma)\csc^2\left[x + \frac{t(2+\gamma)(2+\beta+2\gamma)}{3\theta}\right]}{3\theta}.$$
(3.21)



Figure 3.7: 3D plot of solution (3.21) for $\gamma = 1, \beta = 3, \theta = 1$ values with $-15 \le x \le 15, -4 \le t \le 4$ range and 2D plot of solution for t = 3 with these values.

Case3:

$$a_{0} = \frac{-3\left(\beta - 36(2+\gamma)\right) + \sqrt{3\left(64 + 3\beta^{2} - 16\gamma^{2} - 8\beta(2+\gamma)\right)}}{12\theta}, a_{1} = -\frac{8(2+\gamma)}{\theta}, a_{2} = \frac{2(2+\gamma)}{\theta}, b_{1} = 0, b_{2} = 0,$$

$$\alpha = \frac{\left(3\beta^{2} - 8(2+\gamma)(7+\gamma) + 2\beta(11+7\gamma)\right) + (6-\beta+2\gamma)\sqrt{3\left(64+3\beta^{2} - 16\gamma^{2} - 8\beta(2+\gamma)\right)}}{12\theta},$$

$$v = \frac{4(1+\gamma)(2+\gamma) - \beta(5+\gamma) - (1+\gamma)\sqrt{3\left(64+3\beta^{2} - 16\gamma^{2} - 8\beta(2+\gamma)\right)}}{12\theta}.$$
(3.22)

Embedding the these values in Eqs. (3.5) and (3.17), we acquire the soliton solution of Eq. (1.1)

$$u_{8}(x,t) = \frac{-3\beta + \sqrt{3(64 + 3\beta^{2} - 16\gamma^{2} - 8\beta(2 + \gamma)) + 12(2 + \gamma)(1 + 2\cot^{2}[x - \nu t])}}{12\theta},$$
(3.23)

where $v = \frac{4(1+\gamma)(2+\gamma) - \beta(5+\gamma) - (1+\gamma)\sqrt{3(64+3\beta^2 - 16\gamma^2 - 8\beta(2+\gamma))}}{12\theta}$



Figure 3.8: 3D plot of solution (3.23) for $\gamma = 0.1$, $\beta = -3$, $\theta = 2$ values with $-20 \le x \le 20$, $-2 \le t \le 2$ range and 2D plot of solution for t = 1 with these values.

Case4:

$$a_{0} = \left(-2 + 8\sqrt{11}\right)a_{2}, a_{1} = -2\left(1 + \sqrt{11}\right)a_{2}, b_{1} = -10\left(1 + \sqrt{11}\right)a_{2}, b_{2} = 25a_{2},$$

$$\alpha = -18\left(-4 + \sqrt{11}\right)a_{2}, \beta = 28 - 6\sqrt{11}, \theta = \frac{1}{a_{2}}, \gamma = \frac{-3}{2}, \nu = 6\left(-4 + \sqrt{11}\right)a_{2}.$$
(3.24)

Embedding the these values in Eqs. (3.5) and (3.17), we acquire the soliton solution of Eq. (1.1)

$$u_{9}(x,t) = \frac{\csc^{4}[x-vt]\left(1-4\sqrt{11}+4\left(1+\sqrt{11}\right)\cos[2x-2vt]-3\cos[4x-4vt]-2\left(1+\sqrt{11}\right)\sin[2x-2vt]+4\sin[4x-4vt]\right)a_{2}}{2\left(2+\cot[x-vt]\right)^{2}}, (3.25)$$

where $v = 6(-4 + \sqrt{11})a_2$.



Figure 3.9: 3D plot of solution (3.25) for $a_2 = 4$ values with $-25 \le x \le 25, -3 \le t \le 3$ range and 2D plot of solution for t = 2 with these values.

Family 3: For p = [i, -i, 1, 1] and q = [i, -i, i, -i], Eq. (2.5) turns into the form,

$$\Phi(\eta) = \frac{-\sin(\eta)}{\cos(\eta)}.$$
(3.26)

Case1:

$$a_{0} = \frac{2(2+\gamma)}{\theta}, a_{1} = 0, a_{2} = \frac{2(2+\gamma)}{\theta}, b_{1} = 0, b_{2} = 0,$$

$$\alpha = -\frac{5(2+\gamma)(2+\beta+2\gamma)}{3\theta}, v = -\frac{(2+\gamma)(2+\beta+2\gamma)}{3\theta}.$$
(3.27)

Embedding the these values in Eqs. (3.5) and (3.26), we acquire the soliton solution of Eq. (1.1)

$$u_{10}(x,t) = \frac{2(2+\gamma)\sec^{2}\left[x + \frac{t(2+\gamma)(2+\beta+2\gamma)}{3\theta}\right]}{\theta}.$$
(3.28)



Figure 3.10: 3D plot of solution (3.28) for $\gamma = -0.1$, $\beta = 0.3$, $\theta = 1$ values with $-17 \le x \le 17$, $-2 \le t \le 2$ range and 2D plot of solution for t = 1 with these values.

Case2:

$$\alpha_{0} = \frac{-3\beta\theta + \sqrt{3\left(3\beta^{2} - 32\left(\beta + 8\left(-2 + \gamma\right)\right)\left(2 + \gamma\right)\right)}}{12\theta}, a_{1} = 0, a_{2} = \frac{2(2 + \gamma)}{\theta}, b_{1} = 0, b_{2} = \frac{2(2 + \gamma)}{\theta}, a_{1} = 0, a_{2} = \frac{2(2 + \gamma)}{\theta}, a_{2} = \frac{2(2 + \gamma)}{\theta}, a_{2} = \frac{2(2 + \gamma)}{\theta}, a_{1} = 0, a_{2} = \frac{2(2 + \gamma)}{\theta}, a_{1} = 0, a_{2} = \frac{2(2 + \gamma)}{\theta}, a_{2} =$$

Embedding the these values in Eqs. (3.5) and (3.26), we acquire the soliton solution of Eq. (1.1)

$$u_{11}(x,t) = \frac{-3\beta + \sqrt{3(3\beta^2 - 32\beta(2+\gamma) - 256(-4+\gamma^2)) + 24(2+\gamma)(1+\cot^4[x-\nu t])\tan^2[x-\nu t]}}{12\theta},$$
(3.30)



Figure 3.11: 3D plot of solution (3.30) for $\gamma = -1$, $\beta = 3$, $\theta = 1$ values with $-20 \le x \le 20$, $-5 \le t \le 5$ range and 2D plot of solution for t = 2 with these values.

where $v = \frac{16(1+\gamma)(2+\gamma) - \beta(5+\gamma) - (1+\gamma)\sqrt{3(3\beta^2 - 32\beta(2+\gamma) - 256(-4+\gamma^2))}}{12\theta}$. Family 4: For p = [1, 1, -1, 1] and q = [1, -1, 1, -1], Eq. (2.5) turns into the form,

$$\Phi(\eta) = \frac{-\cosh(\eta)}{\sinh(\eta)}.$$
(3.31)

Case1:

$$a_{0} = -\frac{2(2+\gamma)}{\theta}, a_{1} = 0, a_{2} = \frac{2(2+\gamma)}{\theta}, b_{1} = 0, b_{2} = 0,$$

$$\alpha = \frac{5(2+\gamma)(\beta - 2(1+\gamma))}{3\theta}, v = -\frac{(2+\gamma)(2-\beta+2\gamma)}{3\theta}.$$
(3.32)

Embedding the these values in Eqs. (3.5) and (3.31), we acquire the soliton solution of Eq. (1.1)

$$u_{12}(x,t) = \frac{2(2+\gamma)\operatorname{csch}^2\left[x - \frac{t(2+\gamma)(2-\beta+2\gamma)}{3\theta}\right]}{\theta}.$$
(3.33)



Figure 3.12: 3D plot of solution (3.33) for $\gamma = 3$, $\beta = 0.5$, $\theta = 3$ values with $-10 \le x \le 10$, $-2 \le t \le 2$ range and 2D plot of solution for t = 1 with these values.

Case2:

$$a_{0} = -\frac{3(8+\beta+4\gamma) + \sqrt{3(64+3\beta^{2}-16\gamma^{2}-8\beta(2+\gamma))}}{12\theta}, a_{1} = 0, a_{2} = \frac{2(2+\gamma)}{\theta}, b_{1} = 0, b_{2} = 0,$$

$$\alpha = \frac{3(\beta^{2}-8(2+\gamma)(3+\gamma)-2\beta(7+3\gamma)) + (10+\beta+6\gamma)\sqrt{3(64+3\beta^{2}-16\gamma^{2}-8\beta(2+\gamma))}}{24\theta},$$

$$\nu = \frac{-4(1+\gamma)(2+\gamma) - \beta(5+\gamma) - (1+\gamma)\sqrt{3(64+3\beta^{2}-16\gamma^{2}+8\beta(2+\gamma))}}{12\theta}.$$
(3.34)

Embedding the these values in Eqs. (3.5) and (3.31), we acquire the soliton solution of Eq. (1.1)

$$u_{13}(x,t) = -\frac{3(8+\beta+4\gamma) + \sqrt{3(64+3\beta^2 - 16\gamma^2 + 8\beta(2+\gamma)) - 24(2+\gamma)\coth^2[x-\nu t]}}{12\theta},$$
(3.35)

where $v = \frac{-4(1+\gamma)(2+\gamma) - \beta(5+\gamma) - (1+\gamma)\sqrt{3(64+3\beta^2 - 16\gamma^2 + 8\beta(2+\gamma))}}{12\theta}$



Figure 3.13: 3D plot of solution (3.35) for $\gamma = 3$, $\beta = 0.7$, $\theta = 1$ values with $-30 \le x \le 30$, $-3 \le t \le 3$ range and 2D plot of solution for t = 2 with these values.

Case3:

$$a_{0} = -\frac{3(8+\beta+4\gamma) + \sqrt{3(64+3\beta^{2}-16\gamma^{2}+8\beta(2+\gamma))}}{12\theta}, a_{1} = 0, a_{2} = 0, b_{1} = 0, b_{2} = \frac{2(2+\gamma)}{\theta},$$

$$\alpha = \frac{3(\beta^{2}-8(2+\gamma)(3+\gamma)-2\beta(7+3\gamma)) + (10+\beta+6\gamma)\sqrt{3(64+3\beta^{2}-16\gamma^{2}+8\beta(2+\gamma))}}{24\theta},$$

$$v = \frac{-4(1+\gamma)(2+\gamma) - \beta(5+\gamma) - (1+\gamma)\sqrt{3(64+3\beta^{2}-16\gamma^{2}+8\beta(2+\gamma))}}{12\theta}.$$
(3.36)

Embedding the these values in Eqs. (3.5) and (3.31), we acquire the soliton solution of Eq. (1.1)

$$u_{14}(x,t) = -\frac{3(8+\beta+4\gamma) + \sqrt{3(64+3\beta^2 - 16\gamma^2 + 8\beta(2+\gamma)) - 24(2+\gamma)\tanh^2[x-\nu t]}}{12\theta},$$
(3.37)

where $v = \frac{-4(1+\gamma)(2+\gamma)-\beta(5+\gamma)-(1+\gamma)\sqrt{3(64+3\beta^2-16\gamma^2+8\beta(2+\gamma))}}{12\theta}$.



Figure 3.14: 3D plot of solution (3.37) for $\gamma = -1$, $\beta = -3$, $\theta = 0.25$ values with $-25 \le x \le 25$, $-4 \le t \le 4$ range and 2D plot of solution for t = 3 with these values.

Case4:

$$\alpha = -\frac{3\beta\theta + \sqrt{3\left(3\beta^2 + 32\left(16 + \beta - 8\gamma\right)(2 + \gamma)\right)}}{12\theta}, a_1 = 0, a_2 = \frac{2(2 + \gamma)}{\theta}, b_1 = 0, b_2 = \frac{2(2 + \gamma)}{\theta}, a_1 = 0, a_2 = \frac{2(2 + \gamma)}{\theta}, b_1 = 0, b_2 = \frac{2(2 + \gamma)}{\theta}, a_1 = 0, a_2 = \frac{2(2 + \gamma)}{\theta}, a_2 = \frac{2(2 + \gamma)}{\theta}, a_1 = 0, a_2 = \frac{2(2 + \gamma)}{\theta}, a_2 = \frac{2(2 + \gamma)}{\theta}, a_2 = \frac{2(2 + \gamma)}{\theta}, a_1 = 0, a_2 = \frac{2(2 + \gamma)}{\theta}, a_2 = \frac{2(2 + \gamma)}{\theta}, a_2 = \frac{2(2 + \gamma)}{\theta}, a_1 = 0, a_2 = \frac{2(2 + \gamma)}{\theta}, a_1 = \frac{2(2 + \gamma)}{\theta}, a_2 = \frac{2(2 + \gamma)}{\theta},$$

Embedding the these values in Eqs. (3.5) and (3.31), we acquire the soliton solution of Eq. (1.1)

$$u_{15}(x,t) = -\frac{3\beta + \sqrt{3(3\beta^2 + 32\beta(2+\gamma) - 256(-4+\gamma^2)) - 24(2+\gamma)(1+\coth^4[x-vt])\tanh^2[x-vt]}}{12\theta},$$
(3.39)

where $v = \frac{-16(1+\gamma)(2+\gamma) - \beta(5+\gamma) - (1+\gamma)\sqrt{3(3\beta^2 + 32\beta(2+\gamma) - 256(-4+\gamma^2))}}{12\theta}$.



Figure 3.15: 3D plot of solution (3.39) for $\gamma = -4$, $\beta = 4$, $\theta = 1$ values with $-50 \le x \le 50$, $-5 \le t \le 5$ range and 2D plot of solution for t = 3 with these values.

4. Result and Discussion

Some soliton solutions of the generalized hyperelastic rod wave equation have been obtained by applying GERFM. These results, obtained with the help of Wolfram Mathematica 12, have been graphically represented and their accuracy has been proven. When the previous results related to this equation are compared with the results obtained in this study, our (3.35) and (3.37) solutions are similar to the (42) solutions given by Gozükızıl and Akçagıl. Other solutions we have obtained according to our research have not been shown before and are new.

5. Conclusion

In this study, the generalized hyperelastic-rod wave equation was examined. GERFM, which is the solution method of NLEEs, was applied to this equation and thus some trigonometric function, hyperbolic function, complex hyperbolic function, combo soliton, singular soliton, dark soliton and bright soliton solutions of the equation were obtained. In addition, certain values and ranges were given to the obtained solutions and 2D and 3D graphs of these solutions were drawn with the help of Wolfram Mathematica. The most important advantage of the method used in this study is that a wide variety of solution families can be created. It is a more general method compared to other methods, as it offers a wide variety of solution families.

Article Information

Acknowledgements: The authors would like to express their sincere thanks to the editor and the anonymous reviewers for their helpful comments and suggestions.

Author's contributions: All authors contributed equally to the writing of this paper. All authors read and approved the final manuscript.

Conflict of Interest Disclosure: No potential conflict of interest was declared by the author.

Copyright Statement: Authors own the copyright of their work published in the journal and their work is published under the CC BY-NC 4.0 license.

Supporting/Supporting Organizations: No grants were received from any public, private or non-profit organizations for this research.

Ethical Approval and Participant Consent: It is declared that during the preparation process of this study, scientific and ethical principles were followed and all the studies benefited from are stated in the bibliography.

Plagiarism Statement: This article was scanned by the plagiarism program. No plagiarism detected.

Availability of data and materials: Not applicable.

References

- [1] M.M.A. Khater and A.M. Alabdali, Multiple Novels and Accurate Traveling Wave and Numerical Solutions of the (2+1) Dimensional Fisher-Kolmogorov-Petrovskii-Piskunov Equation, Mathematics, 9(12) (2021), 1-13.
- [2] F. Dusunceli, E. Celik, M. Askin and H. Bulut, New exact solutions for the doubly dispersive equation using the improved Bernoulli sub-equation function method, Indian Journal of Physics, 95(2) (2021), 309-314. [3] G. Bakıcıerler, S. Alfaqeih and E. Mısırlı, Analytic solutions of a (2+1)-dimensional nonlinear Heisenberg ferromagnetic spin chain equation, Physica
- A: Statistical Mechanics and its Applications, 582 (2021), 126255. [4] S. Malik, S. Kumar, K.S. Nisar and C.A. Saleel, Different analytical approaches for finding novel optical solitons withgeneralized third-order nonlinear
- Schrödinger equation, Results in Physics, 29 (2021), 104755.
- [5] T. Aktürk, Y. Gurefe and Y. Pandir, An application of the new function method to the Zhiber–Shabat equation, An International Journal of Optimization and Control: Theories Applications (IJOCTA), 7(3) (2017), 271–274.
- [6] A. Akbulut, M. Kaplan and F. Tascan, The investigation of exact solutions of nonlinear partial differential equations by using exp $(-\phi(\xi))$ method, Optik, 132 (2017), 382-387.
- M. Ünal and M. Ekici, The Double (G'/G, 1/G)-Expansion Method and Its Applications for Some Nonlinear Partial Differential Equations, Journal of [7]
- the Institute of Science and Technology, 11(1) (2021), 599-608. R.I. Nuruddeen, K. S. Aboodh and K. K. Ali, Analytical Investigation of Soliton Solutions to Three Quantum Zakharov-Kuznetsov Equations, Communications in Theoretical Physics, 70(4) (2018), 405-412. [8] M. Tahir and A. U. Awan, Optical singular and dark solitons with Biswas-Arshed model by modified simple equation method, Optik, 202 (2020), [9]
- [10] O. Tasbozan, Y. Çenesiz and A. Kurt, New solutions for conformable fractional Boussinesq and combined KdV-mKdV equations using Jacobi elliptic
- function expansion method, The European Physical Journal Plus, 131 (2016), 244
- [11] H.F. Ismael, H. Bulut and H.M. Baskonus, W-shaped surfaces to the nematic liquid crystals with three nonlinearity laws, Soft Computing, 25 (2021), 4513–4524. [12] H.F. Ismael, M.A.S Murad and H. Bulut, Various exact wave solutions for KdV equation with time-variable coefficients, Journal of Ocean Engineering
- and Science, 7(5) (2022), 409–418. [13] H.F. Ismael and H. Bulut, On the Wave Solutions of (2+1)-Dimensional Time-Fractional Zoomeron Equation, Konuralp Journal of Mathematics, 8(2)
- 2020), 410-418,
- [14] H.F. Ismael, S.S. Atas, H. Bulut and M.S. Osman, Analytical solutions to the M-derivative resonant Davey-Stewartson equations, Modern Physics Letters B, 35(30) (2021), 2150455.
- [15] H.F. Ismael and H. Bulut, Nonlinear dynamics of (2+1)-dimensional Bogoyavlenskii-Schieff equation arising in plasma physics, Mathematical Methods in the Applied Sciences, 44(13) (2021), 10321-10330.
- S. Akcagil, T. Aydemir and O.F. Gozukizil, Exact travelling wave solutions of nonlinear pseudoparabolic equations by using the (G'/G) Expansion [16] Method, New Trends in Mathematical Sciences, 4(4) (2016), 51-66.
- O.F. Gozukizil and S. Akcagil, The tanh-coth method for some nonlinear pseudoparabolic equations with exact solutions, Advances in Difference [17] Equations, 2013 (2013), 143.
- [18] H. Holden and X. Raynaud, Global conservative solutions of the generalized hyperelastic-rod wave equation, Journal of Differential Equations, 233 (2007), 448-484.
- M. Bendahmane, G. Coclite and K. Karlsen, H¹-perturbations of smooth solutions for a weakly dissipative hyperelastic-rod wave equation, Mediterranean [19] Journal of Mathematics, 3 (2006), 419-432.
- [20] G.M. Coclite, H. Holden and K.H. Karlsen, Global Weak Solutions To A Generalized Hyperelastic-Rod Wave Equation, SIAM Journal on Mathematical Analysis, 37(4) (2005), 51-66.
- [21] H.M. Srivastava, H. Günerhan and B. Ghanbari, Exact traveling wave solutions for resonance nonlinear Schrödinger equation with intermodal dispersions and the Kerr law nonlinearity, Mathematical Methods in the Applied Sciences, 42(18) (2019), 7210-7221
- [22] B. Ghanbari, M. S. Osman and D. Baleanu, Generalized exponential rational function method for extended Zakharov Kuzetsov equation with conformable derivative, Modern Physics Letters A, 34(20) (2019), 1950155.
- [23] B. Ghanbari and J.F. Gomez, The generalized exponential rational function method for Radhakrishnan-Kundu-Lakshmanan equation with β -conformable time derivative, Revista Mexicana de Fisica, 65(5) (2019), 503-518. Y. Sağlam Özkan, The generalized exponential rational function and Elzaki-Adomian decomposition method for the Heisenberg ferromagnetic spin [24]
- chain equation, Journal of Nonlinear Science and Applications, 35(12) (2021), 2150200.
- [25] K.K. Ali, R. Yilmazer, H. Bulut, T. Aktürk and M.S. Osman, Abundant exact solutions to the strain wave equation in micro-structured solids, Modern Physics Letters B, 35 (2021), 2150439.