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Jacobi Elliptic Function Solutions of Space-Time Fractional Symmetric Regularized Long Wave Equation

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Abstract

In this paper, by using a direct method based on the Jacobi elliptic functions, the exact solutions of the space-time fractional symmetric regularized long wave (SRLW) equation have been obtained. The elliptic function solutions of a nonlinear ordinary differential (auxiliary) equation $(dF/d\xi)^2 = PF^4(\xi) + QF^2(\xi) + R$ have also been examined. Besides, the solutions have been found in general form including rational, trigonometric and hyperbolic functions. Moreover, the complex valued solutions, periodic solutions, and soliton solutions, have also been gained. Some solutions have been illustrated by the graphics.

Keywords: SRLW equation; Jacobi elliptic function; Nonlinear differential equation; Fractional partial differential equation. *AMS Subject Classification (2020):* Primary: 33E05, Secondary: 35G20; 35L05; 35R11.

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1. Introduction

The symmetric regularized long wave (SRLW) equation is in the form

$$u_{tt} + u_{xx} + uu_{xt} + u_x u_t + u_{xxtt} = 0$$

was first described by Seyler and Fenstermacher [1] in 1984 as a model of the weakly nonlinear ion acoustic and space-charge waves. This equation emerges in various physical applications, such as solitary waves with shallow water waves, ion-acoustic waves in plasma and shallow water waves [2]. The solutions of the SRLW equation has been found by the finite difference method [3], exp-function method [4], (G'/G)-expansion method [5], simplest equation method [6], conservative Crank-Nicolson finite difference scheme [7], generalized Jacobi elliptic function method [8], and different version of finite difference method [9].

In recent years, investigation of the exact solutions for fractional differential equations has been popular in the study of scientific research. An important one of these equations is the fractional SRLW equation. So far, the solutions of the space-time fractional SRLW equation has been investigated by utilizing the sub-equation



method [10], functional variable method [11], exp-function method [11], (G'/G)-expansion method [11], tanh-coth method [2], tan-cot method [2], sech-csch method [2] and sec-csc method [2], a novel (G'/G)-expansion method [12], Riccati equation method [13], rational (G'/G)-expansion method [14], improved *F*-expansion method [15], the extended Jacobi elliptic function expansion method [16], the auxiliary equation method [17], new extended direct algebraic method [18], improved Bernoulli sub-equation function method [19], modified extended tanh method [20], rational $exp(-\Omega(\eta))$ -expansion method [21], (G'/G, 1/G)-expansion method [22], extended auxiliary equation mapping method [23], $(D^{\alpha}G/G)$ -expansion method [24], modified Kudryashov method [25], and the fractional $(D_{\xi}^{\alpha}G/G)$ -expansion method [26]. Among these methods, rational (G'/G)-expansion, new extended direct algebraic, improved Bernoulli sub-equation function, and modified extended tanh methods include the conformable derivatives.

The aim of this paper is to obtain the largest set of exact solutions in the literature of space-time fractional SRLW equation in the form

$$D_t^{\alpha} D_t^{\alpha} u + D_x^{\beta} D_x^{\beta} u + u D_t^{\alpha} \left(D_x^{\beta} u \right) + \left(D_t^{\alpha} u \right) \left(D_x^{\beta} u \right) + D_t^{\alpha} D_t^{\alpha} \left(D_x^{\beta} D_x^{\beta} u \right) = 0,$$

$$(1.1)$$

where $0 < \alpha, \beta < 1$, D_t^{α} and D_x^{β} mean conformable fractional derivative of function u(x, t) with respect to t and x, respectively.

2. Preliminaries

Twelve Jacobi elliptic functions are available in the literature. Basic Jacobi elliptic functions are expressed as

 $sn\xi = sn(\xi;m) = sn(\xi|m^2),$ $cn\xi = cn(\xi;m) = cn(\xi|m^2),$ $dn\xi = dn(\xi;m) = dn(\xi|m^2)$

where *m* is the modulus and is a complex number. When the *m* is real, it can always be arranged $0 < m^2 < 1$. In addition to these functions, Glaisher found the other elliptic functions sd, cd, nd, sc, nc, dc, ns, cs, and ds by taking reciprocals and quotients of basic Jacobi elliptic functions [27]. Besides, when m = 0 and m = 1, Jacobi elliptic functions turn into trigonometric and hyperbolic functions [28].

In recent years, Khalil et al. [29] defined a new fractional derivative which is called conformable fractional derivative. This definition is the simplest of other fractional derivatives because it is similar to the definition of the usual derivative. Therefore, the space-time SRLW equation is considered in conformable sense. The definition and the properties of the conformable derivative are given below.

Definition 2.1. [29] Let $f : [0, \infty) \to R$ be a function. The α -th order conformable fractional derivative of f is defined by

$$D^{\alpha}(f)(t) = \lim_{\epsilon \to 0} \frac{f(t + \epsilon t^{1-\alpha}) - f(t)}{\epsilon}, \qquad t > 0, \qquad \alpha \in (0, 1).$$

If *f* is α -differentiable in some $(0, \alpha)$, $\alpha > 0$ and $\lim_{t \to 0^+} f^{(\alpha)}(t)$ exists, then we define $f^{(\alpha)}(0) = \lim_{t \to 0^+} f^{(\alpha)}(t)$.

Theorem 2.1. [29] Let $\alpha \in (0, 1]$ and suppose f, g are α -differentiable at point t > 0. Then, the following are satisfied:

- 1. $D^{\alpha}(cf + dg) = cD^{\alpha}(f) + dD^{\alpha}(g) \ \forall c, d \in \mathbb{R}.$
- 2. $D^{\alpha}(t^p) = pt^{p-\alpha} \ \forall p \in R.$
- 3. $D^{\alpha}(\lambda) = 0$ for all constant functions $f(t) = \lambda$.
- 4. $D^{\alpha}(fg) = fD^{\alpha}(g) + gD^{\alpha}(f).$
- 5. $D^{\alpha}(\frac{f}{g}) = \frac{gD^{\alpha}(f) fD^{\alpha}(g)}{g^2}.$
- 6. If, in addition, f is differentiable, then $D^{\alpha}(f)(t) = t^{1-\alpha} \frac{df}{dt}$.

Theorem 2.2. [30] Assume $f, g: (0, \infty) \to R$ be α -differentiable functions, where $0 < \alpha \leq 1$. Let h(t) = f(g(t)). Then h(t) is α -differentiable for all t with $t \neq 0$ and $g(t) \neq 0$ we have

$$D^{\alpha}(h)(t) = D^{\alpha}(f)(g(t)).D^{\alpha}(g)(t).g(t)^{\alpha-1}.$$

If t = 0 we have

$$D^{\alpha}(h)(0) = \lim_{t \to 0} D^{\alpha}(f)(g(t)) \cdot D^{\alpha}(g)(t) \cdot g(t)^{\alpha - 1}.$$

3. Solutions of the space-time fractional SRLW eqquation

In this section, utilizing a transformation

$$\xi = k \frac{t^{\alpha}}{\alpha} + l \frac{x^{\beta}}{\beta}$$

such that k and l are constants and utilizing the chain rule, space-time fractional SRLW Eq. (1.1) turns into

$$k^{2}l^{2}\frac{d^{4}u}{d\xi^{4}} + \left(k^{2} + l^{2}\right)\frac{d^{2}u}{d\xi^{2}} + klu\frac{d^{2}u}{d\xi^{2}} + kl\left(\frac{du}{d\xi}\right)^{2} = 0.$$
(3.1)

The main idea of proposed method is to obtain the largest exact solutions of Eq. (3.1) in the form

$$u(\xi) = \sum_{j=0}^{N} a_j F^j(\xi).$$

Here, N and a_j are unknown coefficients. $F(\xi)$ is the solution of nonlinear ordinary differential equation

$$(F')^2(\xi) = PF^4(\xi) + QF^2(\xi) + R$$
 (3.2)

where P, Q and R are constants. This equation has emerged as an auxiliary equation for the solution of many partial differential equations. Differentiating Eq. (3.2), we get the Duffing equation as

$$F^{''}\left(\xi\right) = 2PF^{3}\left(\xi\right) + QF\left(\xi\right)$$

which is used as a mathematical model of various physical systems [31]. The authors are investigated the exact solutions of Eq. (3.2). Some of them are seen in Ref. [32–36]. The Jacobi elliptic function solutions of Eq. (3.2) are presented by Table 1.

Table 1. The Jacobi elliptic function solutions of Eq. (3.2).

	Р	Q	R	F
1	m^2	$-(1 + m^2)$	1	$\pm sn\xi$, $\pm cd\xi$
2	1	$-(1 + m^2)$	<i>m</i> ²	$\pm ns\xi$, $\pm dc\xi$
3	$-m^{2}$	$-(1 + m^2)$	-1	$\pm i \operatorname{sn} \xi$, $\pm i \operatorname{cd} \xi$
4	-1	$-(1 + m^2)$	$-m^{2}$	$\pm ins\xi$, $\pm ide\xi$
5	1	$2 - m^2$	$1 - m^2$	$\pm cs\xi, \pm idn\xi$
6	$1 - m^2$	$2 - m^2$	1	$\pm sc\xi, \pm ind\xi$
7	1	2 m ²	m ² 1	<u> </u>
8	$m^2 - 1$	$2 - m^2$	-1	$\pm i s c \xi, \pm n d \xi$
9	$1 - m^2$	$2m^2 - 1$	$-m^{2}$	$\pm \operatorname{nc}\xi, \pm imsd\xi$
10	$-m^{2}$	$2m^2 - 1$	$1 - m^2$	$\pm cn\xi$, $\pm \frac{i}{m}ds\xi$
11	$m^2 - 1$	$2m^2 - 1$	m^2	$\pm inc\xi$, $\pm msd\xi$
12	m ²	$2m^2 - 1$	$m^2 - 1$	$\pm icn\xi$, $\pm \frac{1}{m}ds\xi$
13	$m^4 - m^2$	$2m^2 - 1$	1	$\pm \frac{i}{m} \operatorname{nc} \xi, \pm \operatorname{sd} \xi$
14	1	$2m^2 - 1$	$m^4 - m^2$	$\pm imen\xi$, $\pm ds\xi$

15	$-m^4 + m^2$	$2m^2 - 1$	-1	$\pm \frac{1}{m} \operatorname{nc} \xi$, $\pm i \operatorname{sd} \xi$
16	-1	$2m^2 - 1$	$-m^4 + m^2$	$\pm m \operatorname{cn} \xi$, $\pm i \operatorname{ds} \xi$
17	$\frac{1}{4}$	$\frac{1 + m^2}{2}$	$\frac{(1-m^2)^2}{4}$	$ds\xi \pm cs\xi, -ds\xi \mp cs\xi,$ $i(men\xi \pm dn\xi), -i(men\xi \pm dn\xi)$
18	$\frac{(1-m^2)^2}{4}$	$\frac{1+m^2}{2}$	$\frac{1}{4}$	$\frac{\frac{\mathrm{sn}\xi}{\mathrm{dn}\xi\pm\mathrm{cn}\xi}, \frac{-\mathrm{sn}\xi}{\mathrm{dn}\xi\pm\mathrm{cn}\xi}}{\frac{i}{\mathrm{mcn}\xi\pm\mathrm{dn}\xi}, \frac{-i}{\mathrm{mcn}\xi\pm\mathrm{dn}\xi}}$
19	$-\frac{1}{4}$	$\frac{1+m^2}{2}$	$-\frac{(1-m^2)^2}{4}$	$\begin{split} &i(\mathrm{ds}\xi\pm\mathrm{cs}\xi),\ -i(\mathrm{ds}\xi\pm\mathrm{cs}\xi),\\ &m\mathrm{cn}\xi\pm\mathrm{dn}\xi,\ -m\mathrm{cn}\xi\mp\mathrm{dn}\xi \end{split}$
20	$-\frac{(1-m^2)^2}{4}$	$\frac{1+m^2}{2}$	$-\frac{1}{4}$	$\frac{\frac{i \operatorname{sn}\xi}{\operatorname{dn}\xi \pm \operatorname{cn}\xi}}{\frac{1}{\operatorname{men}\xi \pm \operatorname{dn}\xi}}, \frac{\frac{-i \operatorname{sn}\xi}{\operatorname{dn}\xi \pm \operatorname{cn}\xi}}{\frac{-1}{\operatorname{men}\xi \pm \operatorname{dn}\xi}},$
21	$\frac{1-m^2}{4}$	$\frac{1+m^2}{2}$	$\frac{1-m^2}{4}$	$\begin{split} & \operatorname{nc}\xi \pm \operatorname{sc}\xi, \ -\operatorname{nc}\xi \mp \operatorname{sc}\xi, \\ & i(\operatorname{msd}\xi \pm \operatorname{nd}\xi), \ -i(\operatorname{msd}\xi \pm \operatorname{nd}\xi), \end{split}$
22	$\frac{m^2 - 1}{4}$	$\frac{1 + m^2}{2}$	$\frac{m^2 - 1}{4}$	$\begin{split} i(\operatorname{nc}\xi \pm \operatorname{sc}\xi), & -i(\operatorname{nc}\xi \pm \operatorname{sc}\xi), \\ & \operatorname{msd}\xi \pm \operatorname{nd}\xi, -\operatorname{msd}\xi \mp \operatorname{nd}\xi, \end{split}$
23	$\frac{1}{4}$	$\frac{m^2 - 2}{2}$	$\frac{m^4}{4}$	$\begin{split} & \mathrm{ns}\xi\pm\mathrm{ds}\xi, \ -\mathrm{ns}\xi\mp\mathrm{ds}\xi, \\ & \mathrm{dc}\xi\pm\sqrt{1-m^2}\mathrm{nc}\xi, \ -\mathrm{dc}\xi\mp\sqrt{1-m^2}\mathrm{nc}\xi \end{split}$
24	$\frac{m^4}{4}$	$\frac{m^2 - 2}{2}$	1 4	$\frac{\frac{\mathrm{sn}\xi}{1\pm\mathrm{dn}\xi'}, \frac{-\mathrm{sn}\xi}{1\pm\mathrm{dn}\xi'}}{\frac{\mathrm{cn}\xi}{\mathrm{dn}\xi\pm\sqrt{1-m^2}}, \frac{-\mathrm{cn}\xi}{\mathrm{dn}\xi\pm\sqrt{1-m^2}}}$
25	$-\frac{1}{4}$	$\frac{m^2 - 2}{2}$	$-\frac{m^4}{4}$	$\begin{split} &i(\mathrm{ns}\xi\pm\mathrm{ds}\xi),\ -i(\mathrm{ns}\xi\pm\mathrm{ds}\xi),\\ &i\left(\mathrm{dc}\xi\pm\sqrt{1-m^2}\mathrm{nc}\xi\right),-i\left(\mathrm{dc}\xi\pm\sqrt{1-m^2}\mathrm{nc}\xi\right)\end{split}$
26	$-\frac{m^4}{4}$	$\frac{m^2 - 2}{2}$	$-\frac{1}{4}$	$\frac{\frac{i \sin \zeta}{1 \pm dn \xi'}}{\frac{i \cosh \zeta}{dn \xi \pm \sqrt{1 - m^2}}}, \frac{\frac{-i \sin \zeta}{1 \pm dn \xi'}}{\frac{-i \cos \xi}{dn \xi \pm \sqrt{1 - m^2}}}$
27	$\frac{m^2}{4}$	$\frac{m^2 - 2}{2}$	$\frac{m^2}{4}$	$\operatorname{sn}\xi \pm \operatorname{icn}\xi, -\operatorname{sn}\xi \mp \operatorname{icn}\xi,$ $\operatorname{cd}\xi \pm i\sqrt{1-m^2}\operatorname{sd}\xi, -\operatorname{cd}\xi \mp i\sqrt{1-m^2}\operatorname{sd}\xi$
28	$-\frac{m^2}{4}$	$\frac{m^2 - 2}{2}$	$-\frac{m^2}{4}$	$\begin{array}{c} \mathrm{cn}\xi\pm\mathrm{isn}\xi,\ -\mathrm{cn}\xi\mp\mathrm{isn}\xi,\\ \sqrt{1-m^2}\mathrm{sd}\xi\pm\mathrm{icd}\xi,\ -\sqrt{1-m^2}\mathrm{sd}\xi\mp\mathrm{icd}\xi\end{array}$
29	$\frac{1}{4}$	$\frac{1-2m^2}{2}$	$\frac{1}{4}$	$\begin{split} & \mathrm{ns}\xi\pm\mathrm{cs}\xi, \ -\mathrm{ns}\xi\mp\mathrm{cs}\xi, \\ & \mathrm{msn}\xi\pm\mathrm{idn}\xi, \ -\mathrm{msn}\xi\mp\mathrm{idn}\xi, \\ & \mathrm{dc}\xi\pm\sqrt{1-m^2}\mathrm{sc}\xi, \ -\mathrm{dc}\xi\mp\sqrt{1-m^2}\mathrm{sc}\xi, \\ & \mathrm{mcd}\xi\pm\mathrm{i}\sqrt{1-m^2}\mathrm{nd}\xi, \ -\mathrm{mcd}\xi\mp\mathrm{i}\sqrt{1-m^2}\mathrm{nd}\xi \end{split}$
30	$-\frac{1}{4}$	$\frac{1-2m^2}{2}$	$-\frac{1}{4}$	$\begin{split} &i(\mathrm{ns}\xi\pm\mathrm{cs}\xi),\ -i(\mathrm{ns}\xi\mp\mathrm{cs}\xi),\\ &\mathrm{dn}\xi\pm\mathrm{i}m\mathrm{sn}\xi,\ -\mathrm{dn}\xi\mp\mathrm{i}m\mathrm{sn}\xi,\\ &i(\mathrm{dc}\xi\pm\sqrt{1-m^2}\mathrm{sc}\xi),\ -i(\mathrm{dc}\xi\pm\sqrt{1-m^2}\mathrm{sc}\xi),\\ &\sqrt{1-m^2}\mathrm{nd}\xi\pm\mathrm{i}m\mathrm{cd}\xi,\ -\sqrt{1-m^2}\mathrm{nd}\xi\mp\mathrm{i}m\mathrm{cd}\xi \end{split}$

Balancing the highest order linear term

$$O\left(\frac{d^4u}{d\xi^4}\right) = N + 4$$

and the highest order nonlinear term

$$O\left(u\frac{d^2u}{d\xi^2}\right) = 2N + 2,$$

 ${\cal N}=2$ is obtained. Therefore, the solution of Eq. (3.1) can be given as

$$u(\xi) = \sum_{j=0}^{2} a_j F^j(\xi) = a_0 + a_1 F + a_2 F^2.$$

Differentiating this equation four times and then substituting the derivatives into Eq. (3.1), sixth order polynomial in F is obtained. Setting its coefficients to be zero, the following equations system is gained,

$$8k^{2}l^{2}RQa_{2} + 2(k^{2} + l^{2})Ra_{2} + 2klRa_{0}a_{2} + klRa_{1}^{2} = 0$$

$$k^{2}l^{2}Q^{2}a_{1} + 12k^{2}l^{2}PRa_{1} + (k^{2} + l^{2})Qa_{1} + kla_{0}Qa_{1} + 6klRa_{1}a_{2} = 0$$

$$16k^{2}l^{2}Q^{2}a_{2} + 72k^{2}l^{2}PRa_{2} + 4(k^{2} + l^{2})Qa_{2} + 4klQa_{0}a_{2} + 2klQa_{1}^{2} + 6klRa_{2}^{2} = 0$$

$$20k^{2}l^{2}PQa_{1} + 2(k^{2} + l^{2})Pa_{1} + 2klPa_{0}a_{1} + 9klQa_{1}a_{2} = 0$$

$$120k^{2}l^{2}QPa_{2} + 6(k^{2} + l^{2})Pa_{2} + 6klPa_{0}a_{2} + 3klPa_{1}^{2} + 8klQa_{2}^{2} = 0$$

$$24k^{2}l^{2}P^{2}a_{1} + 12klPa_{1}a_{2} = 0$$

$$120k^{2}l^{2}P^{2}a_{2} + 10klPa_{2}^{2} = 0.$$

Solving this nonlinear system, the unknown coefficients are found

$$a_0 = B + 4QA, \quad a_1 = 0, \quad a_2 = 12PA$$

such that A = -kl, $B = -(k^2 + l^2)/(kl)$. Hence, the solution of the Eq. (3.1) is

$$u = B + 4QA + 12PAF^2. (3.3)$$

Substituting the P, Q and F given in Table 1 into expression (3.3), exact solutions of Eq. (3.1) are gained and also illustrated by Table 2. Besides, the solutions of space-time fractional Eq. (1.1) can be also obtained by taking inverse transformation.

Table 2. The Jacobi elliptic function solutions of Eq.	(3.2).

1.3	$u = B - 4A(1+m^2) + 12Am^2 \mathrm{sn}^2 \xi,$
-,	$u = B - 4A(1+m^2) + 12Am^2 cd^2\xi$
24	$u = B - 4A(1+m^2) + 12Ans^2\xi,$
-, -	$u = B - 4A(1+m^2) + 12Adc^2\xi$
57	$u = B + 4A(2 - m^2) + 12Acs^2\xi,$
2, 7	$u = B + 4A(2 - m^2) - 12A\mathrm{dn}^2\xi$
6.8	$u = B + 4A(2 - m^2) + 12A(1 - m^2)sc^2\xi,$
0,0	$u = B + 4A(2 - m^2) - 12A(1 - m^2) \mathrm{nd}^2 \xi$
9, 11,	$u = B + 4A(2m^2 - 1) + 12A(1 - m^2)nc^2\xi,$
13, 15	$u = B + 4A(2m^2 - 1) + 12A(-m^2 + m^4)sd^2\xi$
10, 12,	$u = B + 4A(2m^2 - 1) - 12Am^2 \mathrm{cn}^2 \xi,$
14, 16	$u = B + 4A(2m^2 - 1) + 12Ads^2\xi$

17, 19	$u = B + 2A(1 + m^2) + 3A(\mathrm{ds}\xi \mp \mathrm{cs}\xi)^2,$ $u = B + 2A(1 + m^2) - 3A(\mathrm{mcn}\xi \mp \mathrm{dn}\xi)^2$
18, 20	$u = B + 2A(1 + m^2) + 3A(1 - m^2)^2 \operatorname{sn}^2 \xi / (\operatorname{dn} \xi \mp \operatorname{cn} \xi)^2,$ $u = B + 2A(1 + m^2) - 3A(1 - m^2)^2 / (\operatorname{mcn} \xi \mp \operatorname{dn} \xi)^2$
21, 22	$u = B + 2A(1 + m^2) + 3A(1 - m^2)(\operatorname{nc}\xi + \operatorname{sc}\xi)^2,$ $u = B + 2A(1 + m^2) - 3A(1 - m^2)(\operatorname{msd}\xi + \operatorname{nd}\xi)^2$
23, 25	$u = B + 2A(m^{2} - 2) + 3A(\operatorname{ns}\xi \mp \operatorname{ds}\xi)^{2},$ $u = B + 2A(m^{2} - 2) + 3A\left(\operatorname{dc}\xi \mp \sqrt{1 - m^{2}}\operatorname{nc}\xi\right)^{2}$
24, 26	$u = B + 2A(m^{2} - 2) + 3Am^{4} \operatorname{sn}^{2} \xi / (1 + \operatorname{dn} \xi)^{2},$ $u = B + 2A(m^{2} - 2) + 3Am^{4} \operatorname{cn}^{2} \xi / \left(\sqrt{1 - m^{2}} + \operatorname{dn} \xi\right)^{2}$
27, 28	$u = B + 2A(m^{2} - 2) + 3Am^{2}(\operatorname{sn}\xi \mp i\operatorname{cn}\xi)^{2},$ $u = B + 2A(m^{2} - 2) + 3Am^{2}\left(\operatorname{cd}\xi \mp i\sqrt{1 - m^{2}}\operatorname{sd}\xi\right)^{2}$
29, 30	$u = B + 2A(1 - 2m^{2}) + 3A(ns\xi \mp cs\xi)^{2},$ $u = B + 2A(1 - 2m^{2}) + 3A(msn\xi \mp idn\xi)^{2},$ $u = B + 2A(1 - 2m^{2}) + 3A\left(dc\xi \mp \sqrt{1 - m^{2}}sc\xi\right)^{2}$ $u = B + 2A(1 - 2m^{2}) + 3A\left(mcd\xi \mp i\sqrt{1 - m^{2}}nd\xi\right)^{2}$

The elementary function solutions of Eq. (3.1) are also obtained in Table 3 by using the solutions in Table 2 and the Jacobi elliptic functions for m = 0 and m = 1.

m = 0	m = 1
u=B-A,	u=B+A,
u=B-4A,	u=B+4A,
$u = B + 8A + 12A \tan^2 \xi,$	$u = B - 8A + 12A \tanh^2 \xi,$
$u = B + 8A + 12A\cot^2\xi,$	$u = B - 8A + 12A \coth^2 \xi,$
$u = B - 4A + 12A\mathrm{sec}^2\xi,$	$u = B + 4A - 12A \mathrm{sech}^2\xi,$
$u = B - 4A + 12A\csc^2\xi,$	$u = B + 4A + 12A \operatorname{csch}^2 \xi,$
$u = B + 2A + 3A(\sec\xi \pm \tan\xi)^2,$	$u = B - 2A + 3A(\tanh\xi \pm i\mathrm{sech}\xi)^2$,
$u = B + 2A + 3A(\csc\xi \pm \cot\xi)^2,$	$u = B - 2A + 3A(\operatorname{csch}\xi \pm \operatorname{coth}\xi)^2,$
$u = B + 2A + 3A\sin^2\xi/(1\pm\cos\xi)^2$	$u = B - 2A + 3A \tanh^2 \xi / (1 \pm \operatorname{sech} \xi)^2$

Table 3: The elementary function solutions of Eq. (3.2).

4. Demonstrations

In this section, three solutions of space-time fractional SRLW equation from Table 2 and Table 3 are given. These solutions are demonstrated by the aid of Mathematica 11. 3.

In all figures, the solutions are investigated for k = l = 1, and utilizing these constants, A = -1 and B = -2. Firstly, let us consider the periodic solution

$$u = -2 - 4(2m^2 - 1) + 12m^2 \mathrm{cn}^2 \xi$$

Figure 1 illustrates this solution for $-5 \le \xi \le 5$ and $0 \le m \le 1$. Besides, Figure 2 demonstrates 2D graph of the same solution for $-5 \le \xi \le 5$ and different *m* values; namely, the line with dots represents the solution when m = 0, the unitary line represents the solution when m = 0.5, the line with dashes represents the solution when m = 0.8 and the line with dotdashes represents the solution when m = 1.

Clearly seen from Figure 2 that the wave amplitudes are constant for any chosen m value but the wave amplitudes and the wavelengths are increasing while *m* goes from 0 to 1.



Figure 1. 3D graph of the solution $u(\xi, m)$ when $0 \le m \le 1$.



Figure 2. 2D graph of the solution $u(\xi, m)$ when m = 0, m = 0.5, m = 0.8, m = 1.

Secondly, let us consider space-time fractional SRLW Eq. (1.1) for $\alpha = \beta = 0.5$; that is

$$D_t^{1/2} D_t^{1/2} u + D_x^{1/2} D_x^{1/2} u + u D_t^{1/2} \left(D_x^{1/2} u \right) + \left(D_t^{1/2} u \right) \left(D_x^{1/2} u \right) + D_t^{1/2} D_t^{1/2} \left(D_x^{1/2} D_x^{1/2} u \right) = 0.$$
(4.1)

Utilizing the transformation $\xi = 2\sqrt{t} + 2\sqrt{x}$, the solution $u = B + 8A + 12A \tan^2 \xi$ in Table 3 turns into

$$u(x,t) = -10 - 12\tan^2(2\sqrt{t} + 2\sqrt{x}).$$

Figure 3 demonstrates this solution for $0 \le x \le 2$ and $0 \le t \le 1$.Besides, Figure 4 illustrates the same solution for $0 \le x \le 100$ and t = 1. Here, the wave amplitudes are goes to infinity, and the wavelengths are increasing when x increases for $0 \le x < \infty$.

Finally, let us consider space-time fractional SRLW Eq. (4.1). Using the transformation $\xi = 2\sqrt{t} + 2\sqrt{x}$, the solution $u = B - 8A + 12A \coth^2 \xi$ in Table 3 turns into

$$u(x,t) = 6 - 12 \coth^2(2\sqrt{t} + 2\sqrt{x}).$$

Figure 5 illustrates this solution for $0 \le x \le 4$ and $0 \le t \le 1$. Moreover, Figure 6 demonstrates the same solution for $0 \le x \le 4$ and t = 1.



Figure 3. 3D graph of the solution u(x, t) when m = 0.



Figure 4. 2D graph of the exact solution u(x, t) when m = 0 and t = 1.



Figure 5. 3D graph of the solution u(x, t) when m = 1.



5. Conclusions

In this paper, a direct method based on the Jacobi elliptic functions is presented to gain the exact solutions of space-time fractional SRLW equation. The suggested method has many advantages such that the solutions are found in the form including the rational, trigonometric, hyperbolic functions. The complex valued solutions, soliton solutions, and periodic solutions are also obtained. Some of these solutions are illustrated by two-dimensional and three-dimensional graphics. The other advantage of the proposed method is not to require perturbation, linearization, initial and boundary conditions. Besides, solutions of an auxiliary nonlinear ordinary differential equation have been investigated, and given by table. Solving this auxiliary equation, solutions of Duffing equation are also found. Moreover, using this equation, many partial differential equations can be solved.

In the literature, rational (G'/G)-expansion [14], new extended direct algebraic [18], improved Bernoulli subequation function [19], and modified extended tanh [20] methods include the conformable derivatives. 3 solutions which are trigonometric and hyperbolic are obtained by rational (G'/G)-expansion method, 37 solutions which are rational, exponential, trigonometric and hyperbolic are obtained by new extended direct algebraic method, 3 solutions which are rational and exponential are obtained by improved Bernoulli sub-equation function method and 12 solutions which are trigonometric, hyperbolic are obtained by modified extended tanh method. Besides, 8 solutions which are trigonometric, hyperbolic and Jacobi elliptic function by the extended Jacobi elliptic function expansion method [16]. When compared with these methods, our method has the largest number of solutions. Because there are 192 type linear independent solutions for 30 different cases in the presented method. The number of the solutions are infinite depending on the parameters P, Q, R, K and m. When solving different differential equations, the solutions given in Table 1 of auxiliary equation (3.2) can also be utilized. Moreover, these solutions can be helpful for different solution methods. Therefore, these solutions contain the widest set of solutions in the literature.

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